



O&P analysis [13]. Even though the pile up effect is well described and predicted by finite element analysis [14, 15], we still rely on adequate imprints observation (AFM or SEM) to account for it [16-18].

Newly developed methods based on the deformation energies which do not rely on the displacement measurement [19-21] can be a solution for avoiding the pile-up effect. However, these methods need to be validated experimentally to get recognition. Oliver and Pharr method with the necessary recommended calibration procedures (thermal drift, initial contact depth, instrument compliance and indenter geometry) and two methods based on deformation energies [19,20] are presented in the following.

### Oliver and Pharr method

A typical load-displacement curve with principal parameters used in Oliver and Pharr's analysis is represented in Figure 1. The definition of the indentation hardness  $H_{IT}$  and the indentation modulus  $E_{IT}$  using the notation specified in the Standard issued for instrumented indentation [21] are given below.

$$H_{IT} = F_{max}/A_p \quad (1)$$

where  $F_{max}$  is the peak load and  $A_p$  is the projected contact area at that load, defined as a function of the displacement  $h_c$ , representing the indenter contact depth:

$$A_p = F(h_c) \quad (2)$$

with

$$h_c = h_{max} - \varepsilon F_{max}/S \quad (3)$$

where  $S$  is the contact stiffness ( $S = dF/dh$ ) corresponding to the slope of the upper portion of the unloading curve and  $\varepsilon$  is a constant depending on the indenter geometry whose value is 0.75 for most used indenters assimilated to a paraboloid of revolution.

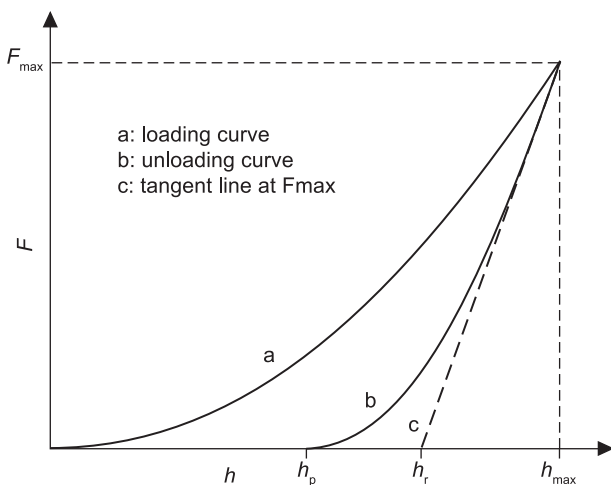


Figure 1. Typical indentation load-displacement curve with principal parameters (from ISO/DIS 14577 norm).

Taking into account the non rigidity of the indenter, the reduced modulus  $E_r$  is related to the contact compliance  $C = dh/dF$  (inverse of the contact stiffness  $S$ ) and the projected contact area  $A_p$  through the following relation:

$$E_r = \frac{\sqrt{\pi}}{2\beta.C\sqrt{A_p}} \quad (4)$$

The dimensionless parameter  $\beta$  is introduced to correct the lack of axial symmetry of pyramidal indenters (deviation of pyramidal indenters from the conical shape). It is usually taken as 1.034 for a Berkovich indenter [22]. In their revised analysis [11] O&P consider that a value of 1.05 is an acceptable choice.

The indentation modulus  $E_{IT}$  of the specimen material is related to the reduced elastic modulus  $E_r$  and the indenter modulus  $E_i$  through the relation:

$$E_{IT} = \frac{1 - \nu_s^2}{\frac{1}{E_r} - \frac{1 - \nu_i^2}{E_i}} \quad (5)$$

where  $\nu_i$  and  $\nu_s$  represent the Poisson coefficient for respectively the indenter and the specimen. The equation (4), derived for elastic contact can be used in elastic-plastic contact for any axisymmetric indenter.

### Calibration procedures

#### Thermal drift

The thermal drift corresponds to changes in the instrument dimensions (expansion or contraction) caused by the surrounding temperature variations. Even if the instrumented indentation equipments are usually well isolated thermally from the environment, small temperature fluctuations can result into errors on the measured indenter displacements, particularly when working at a submicron scale. The thermal drift calibration is usually incorporated in the indentation test procedure. It consists on monitoring the displacement while holding constant the applied load at 10% of the peak load before complete unloading during a short period (~1mn) [10]. The rate of displacement “ $\tau$ ” obtained by linear regression, supposed constant during indentations tests, is used to correct the measured displacement data according to their instant of acquisition. The choice of a small load (10%of the peak load) enables to avoid or reduce any creep effect that can mask the drift effect.

#### Initial penetration depth

With the smallest load allowed by the used indentation instrument according to its resolution limit, a contact between the indenter and the sample is established previously to any subsequent tests. This contact induces an initial penetration depth  $h_i$  that should be added to the measured displacements.

For static loading systems, The ISO/DIS 14577-1 draft standard recommends a polynomial fitting of the data up to 10% of the maximal indentation depth during the loading half cycle considering an elastic-plastic contact. The initial depth could therefore be estimated by extrapolation to zero load [13]. On the other hand, if we consider that the response of the material is elastic during the initial loading (smallest first loading points) with a blunted Berkovich indenter, we can use the following relation:

$$h = k(F^p - F_i^p) \quad (6)$$

The constant  $k$  is also related to the indenter shape and  $F_i$  is the initial contact applied load. The determination of  $p$  and  $k$  enables to evaluate  $h_i$  as  $(k F_i^p)$ .

#### Instrument compliance

The instrument (loading frame) compliance can have an important influence on the measured displacement or the measured compliance, particularly at large depths. The instrument compliance should be established by the manufacturer [23]. However, considering eventual changes in the loading frame components, a simple calibration method is proposed using a reference material (fused silica) [11].

The total measured compliance ( $C = dh/dF$ ) is the sum of the sample compliance  $1/S$  and the instrument compliance. Using equation (4) for defining the sample compliance, the measured compliance is:

$$C = \sqrt{\frac{\pi}{A_p}} \frac{1}{2\beta E_r} + C_f \quad (7)$$

Introducing for the case of an ideal Berkovich indenter surface area ( $A_p = 24.5h_c^2$ ), which is acceptable only for large depth values, the measured compliance becomes:

$$C = \sqrt{\frac{\pi}{24.5}} \frac{1}{2\beta E_r} \frac{1}{h_c} + C_f \quad (8)$$

From this equation, we can see that the plot of  $C$  versus  $1/h_c$  corresponds to a line whose intercept is the instrument compliance  $C_f$ .

This equation can also be rewritten in terms of the peak load  $F_{max}$  instead of the contact load  $h_c$  by replacing  $A_p$  by  $(F_{max} / H_{IT})$  using equation 1 as:

$$C = \frac{\sqrt{\pi}}{2\beta} \frac{1}{E_r} \frac{\sqrt{H_{IT}}}{\sqrt{F_{max}}} + C_f \quad (9)$$

where  $C_f$  is the intercept of the linear plot of  $C$  versus  $1/(F_{max})^{1/2}$ .

In deriving the loading frame compliance by either the equation 8 or 9, it is recommended to discard the low values of  $h_c$  where the indenter geometry can have an

influence on the constancy of the specimen properties  $E_r$  and  $H_{IT}$  and therefore induce errors on the linear regression [13,24].

It is also important to recall that there is a measurable experimental parameter  $F/S^2$  (peak load divided by the stiffness squared) of the tested material that is independent of the penetration depth. It is related to the ratio  $H_{IT}/E_r^2$  according to the following expression deduced from the equations (1) and (4):

$$\frac{F}{S^2} = \frac{\pi}{(2\beta)^2} \frac{H_{IT}}{E_r^2} \quad (10)$$

This parameter, whose value should be constant at high depths where there is no influence of the indenter geometry, is useful in the determination or the verification of the instrument compliance [11].

#### Indenter geometry

The indenter area function  $A_p = F(h_c)$  characterizes the variation of the projected surface area to account for deviation of the indenter geometry from the ideal geometry (the ideal surface area  $A_p$  for a Berkovich indenter is  $24.5h_c^2$ ). This deviation is particularly important at small depths (near the indenter tip) when a blunted or imperfectly machined indenter is used. The area function can be obtained from equation (4) on a reference material (fused silica) knowing its corresponding value  $E_r$  [10].

#### Methods based on the energy of deformation

The elastic, plastic and total indentation work energy can be derived from the load- displacement integrated areas represented in Figure 2.

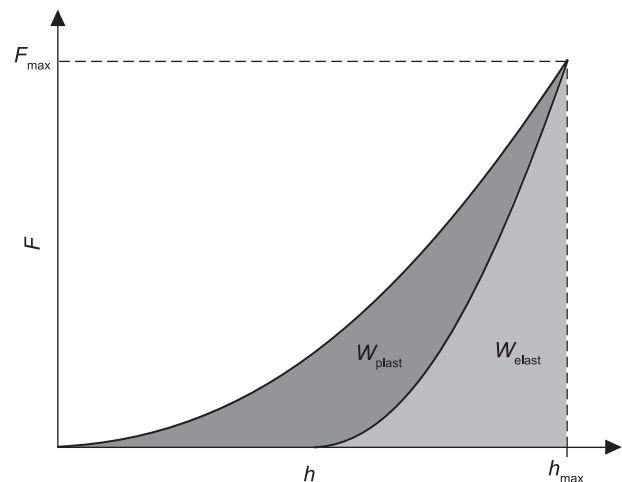


Figure 2. Elastic and plastic indentation work (from ISO/DIS 14577 norm).

Method of Cheng & Cheng

On the basis of a scaling analysis and finite element simulations, Cheng and Cheng (C&C) [19] established a linear correlation between the ratio of the irreversible work to the total work  $(W_t - W_e)/W_t$  and the ratio  $H_{IT}/E_r$  which leads to the following expression in terms of the elastic work over the total work  $W_e/W_t$ :

$$W_e/W_t = \lambda H_{IT}/E_r \tag{11}$$

where  $\lambda$  is a constant that depends on the indenter geometry. For a Berkovich indenter ( $\alpha = 70.3^\circ$ ), different values of  $\lambda$  were proposed in literature [25-28]. The use of relation (11) with the relation (10) defining the parameter  $F/S^2$  enables to determine separately the properties  $H_{IT}$  and  $E_r$ . Such a method which exempts us from the displacement measurements and all the necessary corresponding corrections is surely attractive because the indentation total and elastic works can be more accurately evaluated than the contact depth. It is particularly an interesting way to avoid the problem of an eventual pile up effect.

Method of Chen & Bull

J. Chen and S. J. Bull [20] questioned the linear relationship between  $W_e/W_t$  and  $H_{IT}/E_r$  observing that  $\lambda$  depends on the indenter geometry and also on the tested material. They developed a new relation between  $W_e/W_t$  and  $H_{IT}/E_r$  that is non linear and in which intervene the work hardening exponent  $n$ :

$$\frac{W_e}{W_t} = \frac{1.5(1.24 + 0.2n)\pi}{(0.24 + 0.2n)\beta} \frac{H_{IT}}{E_r} \tan \alpha \left( \frac{1}{1 + \frac{\pi \varepsilon H_{IT}}{2\beta E_r} \tan \alpha} \right) \tag{12}$$

Comparing their relation with other developed relations [15,30,31,32], the authors noticed a better agreement with data. Using relations (1) and (4), they expressed separately the properties  $E_r$  and  $H_{IT}$  as:

$$E_r = \frac{W_e/W_t}{\frac{1.5\pi(1.24 + 0.2n)}{(2.24 + 0.2n)} - \pi \varepsilon W_e/2W_t} \cot \alpha / \frac{4\beta F_{max}}{\pi S^2} \tag{13}$$

$$H_{IT} = \left( \frac{W_e/W_t}{\frac{1.5\pi(1.24 + 0.2n)}{(2.24 + 0.2n)} - \pi \varepsilon W_e/2W_t} \cot \alpha \right)^2 / \frac{4F_{max}}{\pi S^2} \tag{14}$$

Table 1. Glass samples chemical composition.

Components (%)	SiO <sub>2</sub>	Al <sub>2</sub> O <sub>3</sub>	Na <sub>2</sub> O	K <sub>2</sub> O	CaO	MgO	Fe <sub>2</sub> O <sub>3</sub>	MnO	B <sub>2</sub> O <sub>3</sub>	TiO <sub>2</sub>	P <sub>2</sub> O <sub>5</sub>	ZrO <sub>2</sub>	Cr <sub>2</sub> O <sub>3</sub>
SLG	72.4	1.26	13.4	0.24	8.53	3.95	0.16	-	-	0.063	0.018	0.05	//
BSG	80.5	2.49	3.60	0.66	0.21	0.14	0.20	0.002	12.1	0.025	0.016	0.077	0.008

Similarly to the C&C method, the evaluations of the elastic modulus  $E_r$  and the indentation hardness  $H_{IT}$  through equations (13) and (14) do not require any displacement or surface area measurement.

The main objective of this study is to compare the use of displacement O&P method with the two presented energy methods (C&C and C&B methods) in evaluating the mechanical properties  $H_{IT}$  and  $E_r$  of a soda lime glass and a borosilicate glass on a Micro Materials Nano Test. The properties values obtained by conventional means are taken as reference in this comparison. We begin by examining the relative effects of the different experimental sources of errors (initial depth, thermal drift, instrument compliance and indenter geometry) on the derived properties using O&P analysis according to the corresponding implemented calibrations.

EXPERIMENTAL

A soda lime glass (SLG) and a Pyrex borosilicate glass (BSG) with known properties were tested using a Micro Materials Ltd Nano Test equipment (MML NanoTest). The SLG and BSG samples of dimensions 10×10 mm<sup>2</sup> were prepared from glass sheets of respectively 5 mm and 1.1 mm thickness. In order to eliminate any residual stresses on their initial polished surface state, the samples were submitted to an annealing treatment during two hours at 580°C. The chemical compositions of the two glasses, characterized by the X fluorescence and the bolometric methods, are presented on Table 1. Their main characteristics are given in Table 2.

The MML Nano Test equipment used is a static system that functions with a horizontal loading mechanism. It is characterized by a loading range (0-500 mN) with a force and a displacement resolutions of respectively 0.1 mN and 0.1 nm. The equipment is well protected against significant thermal drift, ambient air flow and acoustic disturbance. It is equipped with a Nano-K bench top vibration isolation platform. The nanoindentations tests were done using a Berkovich indenter with an apex semi-angle ( $\Theta = 65.27^\circ$ ) and effective cone angle ( $\alpha = 70.3^\circ$ ). The projected contact area at a contact depth  $h_c$  with a Berkovich indenter corresponds to  $24.5 h_c^2$ .

Table 2. Glass samples characteristics determined conventionally.

	Hardness (GPa)	Elastic modulus (GPa)	Poisson coefficient	Refractive index	Density (g/cm <sup>3</sup> )
SLG	5.5	74	0.23	1.517	2.46
BSG	6.4	63	0.17	1.473	2.23

RESULTS AND DISCUSSION

Prior to the presentation of the experimental results corresponding to the different methods, typical load displacement curves obtained on SLG and BSG samples using multiloading indentations are shown in Figure 3. These reveal continuity on the loading cycles. The dwell times at peak load and at 10 % of the peak load, for respectively the creep and thermal drift effects, are apparent on these curves.

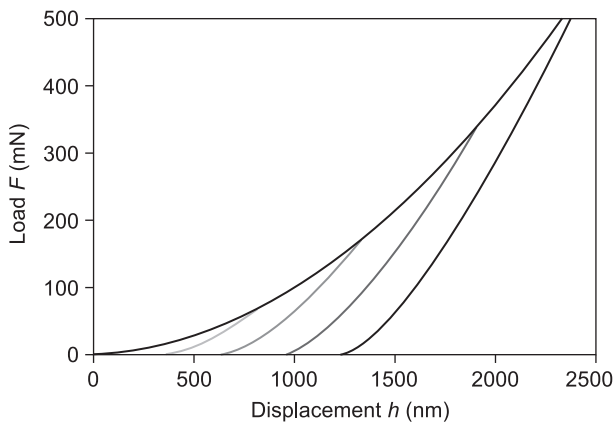
Oliver and Pharr method tests

Initial depth determination

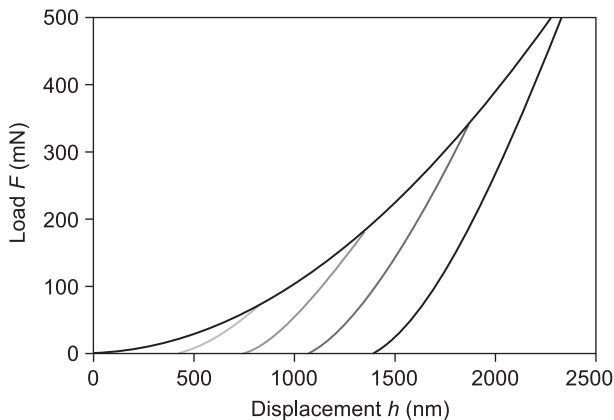
Considering that the contact is elastic for the initial six loading points, we determined the constant  $k$  and  $p$  from equation (6) written in logarithm form as:

$$\log h = \log k + A \log (F^p - F_i^p) \quad (15)$$

We used for that purpose an iterative procedure starting with the hertzian contact value  $2/3$  for  $p$  [13] and with the initial used load ( $F_i = 0.03$  mN). Figure 4



a) BSG



b) SLG

Figure 3. Examples of multiloading indentation curves on glass samples.

represents the plot of two curves obtained on the two glasses after convergence where the slope  $A$  approaches unity for a value of  $p = 0.68$ . The corresponding constant  $k$  value appears to be different for the glasses (63.87 and 54.00) for the SLG and BSG respectively). Using the relation ( $h_i = k F_i^p$ ), the initial obtained depths  $h_i$  are respectively 5.88 nm and 4.97 nm, values to be added to the experimental depths. These initial depths were found to have a limited influence on the determined properties  $H_{IT}$  and  $E_r$ . A variation of less than 1 % and 0.5 % was observed on  $H_{IT}$  and  $E_r$  in general.

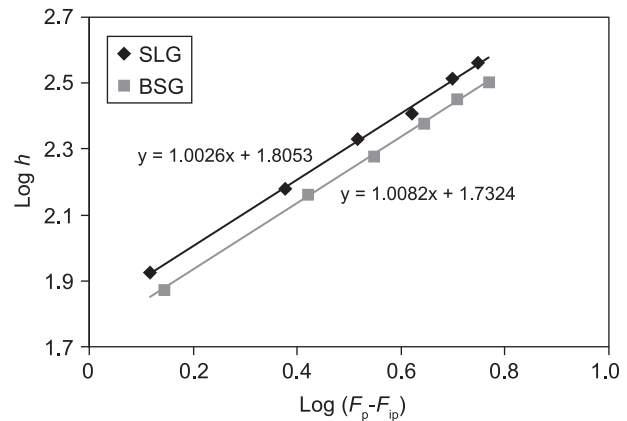


Figure 4. Plots of  $h$  versus  $(F^p - F_i^p)$  for the two glass samples used for determining the initial depth.

Effect of thermal drift

Before considering the thermal drift effect on the determined experimental depths, we remind that the peak load is retained for a period during 30 seconds before unloading in every test in order to reduce significantly the creep effect on the unloading process. The creep value  $C_{IT}$  for both glasses, corresponding to the relative change of the indentation depth as a percentage obtained from Figure 5, does not exceed 2.5%.

Examples of thermal drift curve for both glass samples are shown on Figure 6. They were obtained from a 60 seconds dwell at 10% of the peak load during the unloading half cycle. The thermal drift rate obtained on the linear part (last 40 seconds) corresponds to ( $\tau = -0.145$  nm/s). The experimental depths obtained should be corrected according to their acquisition time  $t$  by adding ( $\tau t$ ). These corrections revealed their relative importance at large depths as we can observe from the changes on a typical curve obtained on a SLG sample using a peak load of 500 mN on Figure 7 representing the separate effect of the thermal drift. Even though, the thermal drift effect seems more important than the initial depth effect, the variations caused by it do not exceed 2 % and 1 % respectively on the measured  $H_{IT}$  and  $E_r$  properties at large depths.

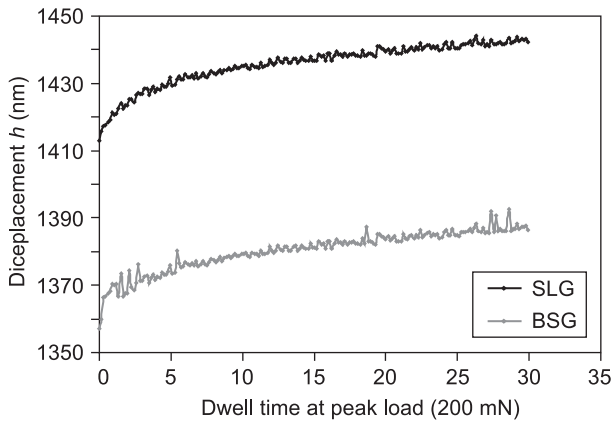


Figure 5. Creep effect during dwell time at  $F_{max} = 200$  mN for SLG and BSG samples.

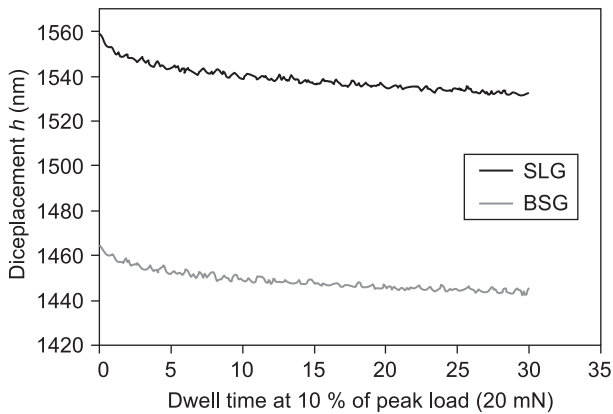
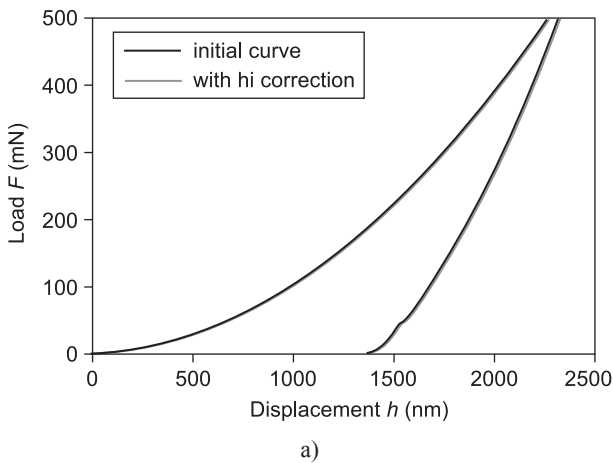


Figure 6. Thermal drift effect on SLG and BSG samples at 10% of  $F_{max} = 200$  mN.



### Instrument compliance effect

We evaluated the instrument compliance by linear regression by plotting  $dh/dF$  versus  $1/h_c$  according to equation 8 and also  $dh/dF$  versus  $1/\sqrt{F_{max}}$  according to equation 9 for each tested glass. The instrument compliance corresponds to the intercept of each line obtained from these plots. As shown in Figure 8, the instrument compliance  $C_f$  values obtained from the first plot (0.0481 and 0.0962 for respectively the BSG and the SLG samples) are much larger than the values obtained from the second plot (0.0138 and 0.0007). These values were obtained from multiloading indentations with peak loads varying between 80 mN and 500 mN and inducing contact depths between 700 nm and 1900 nm. Small depth values were discarded as it is recommended to avoid significant error on the determined slopes due to the indenter geometry [13].

To verify the obtained instrument compliance, we evaluated the parameter  $F_{max}/S^2$  variation with the contact depth  $h_c$  according to equation 10 using the instrument compliance  $C_f$  values obtained through the preceding plots as shown in Figure 9. This figure shows that the constancy of the parameter  $F_{max}/S^2$  is assured only when the instrument compliance contribution approaches zero. The fact that  $C_f$  was found negligible means the correction due to the instrument compliance as established in the last calibration ( $C_f=0.4633$ ) and introduced in the analysis of the instrument is still valid in our case. Without this instrument compliance calibration, errors on the properties measured  $H_{IT}$  can reach 15 %.

### Effect of indenter geometry

Determined by calibration on fused silica, the area function  $A_p(h_c)$  introduced in the analysis after a poly-

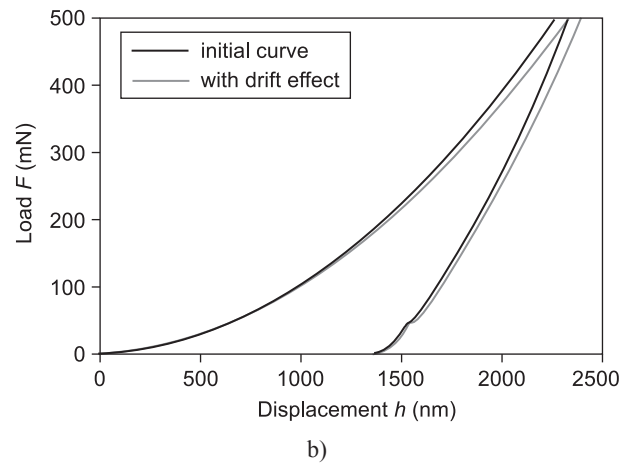


Figure 7. Effect of the initial depth (a) and the thermal drift (b) on a SLG load displacement curve obtained with a peak load ( $F_{max} = 500$  mN).



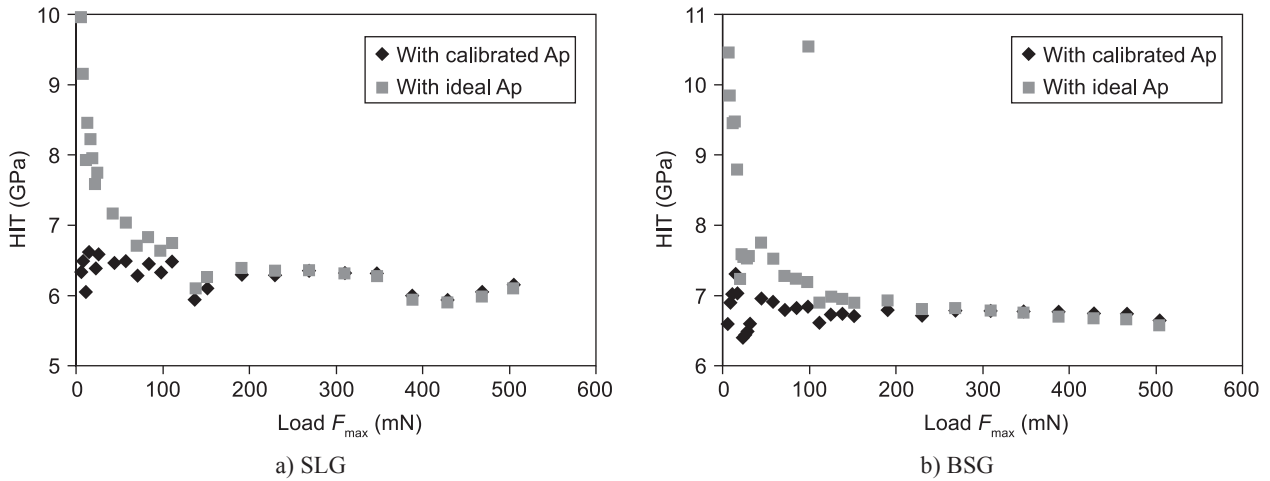


Figure 10. Effect of the calibrated surface area on the indentation hardness.

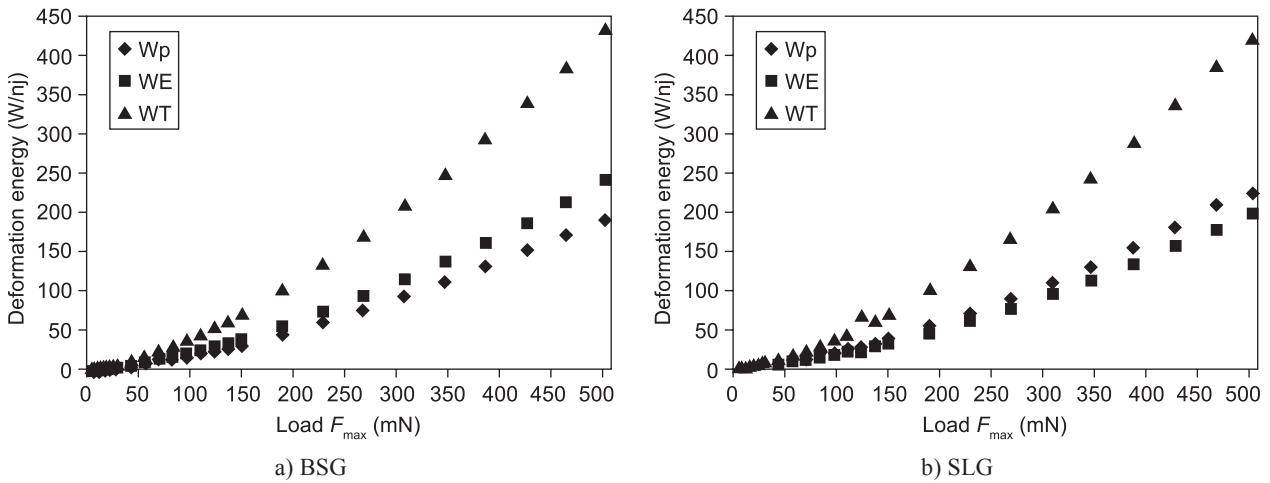


Figure 11. Variation of the elastic, plastic and total indentation energy with peak load  $F_{max}$ .

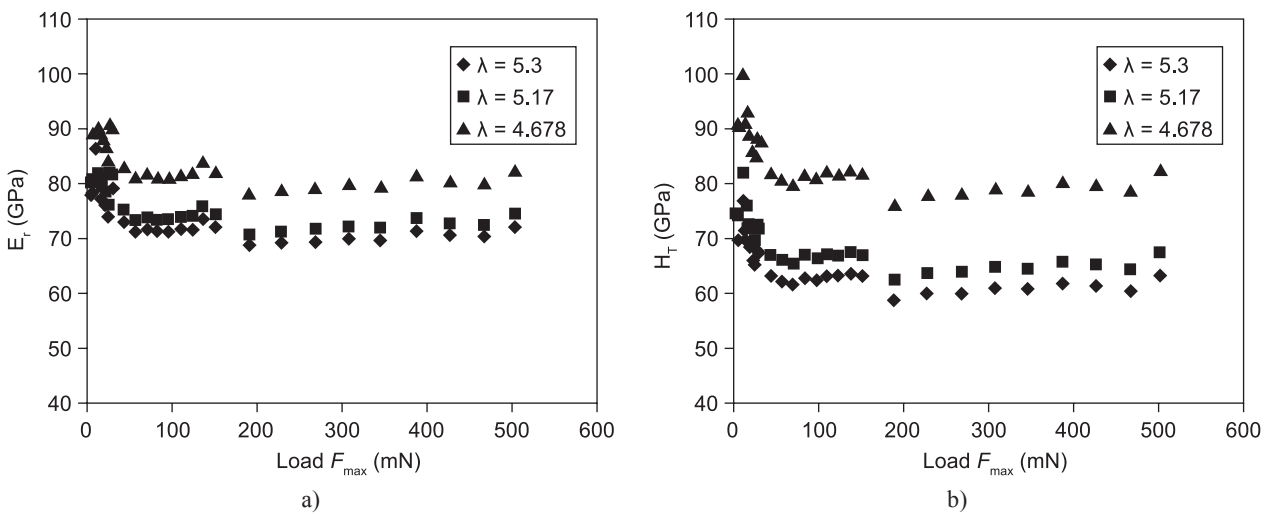


Figure 12. Influence of the constant  $\lambda$  on  $E_r$  (a) and  $H_T$  (b) obtained by Cheng and Cheng method [19] on SLG.



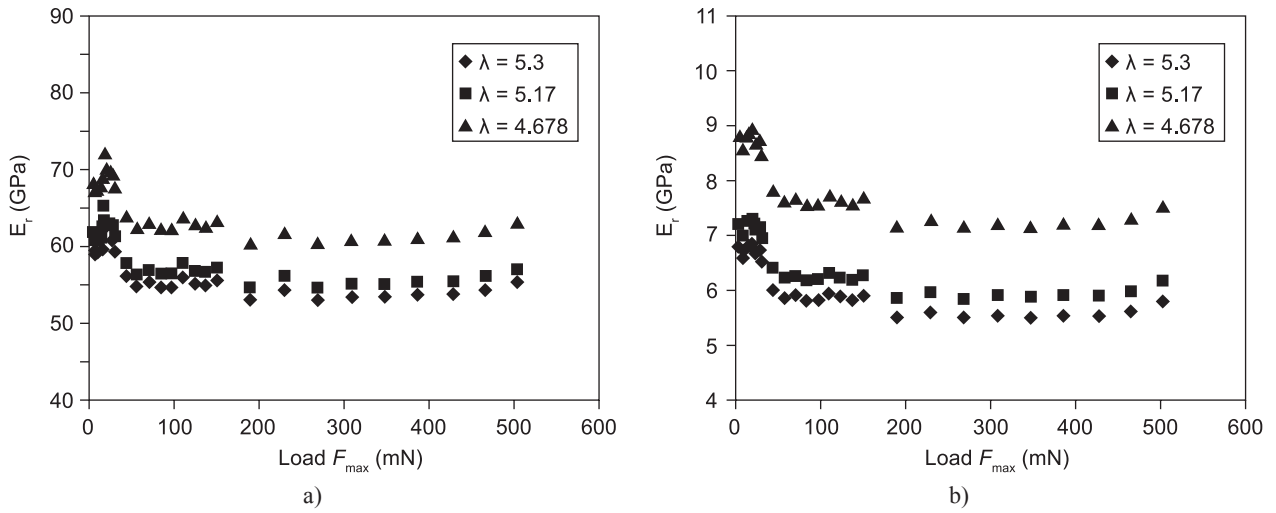


Figure 13. Influence of the constant  $\lambda$  on  $E_r$  (a) and  $H_{ITr}$  (b) obtained by Cheng & Cheng method [19] on BSG.

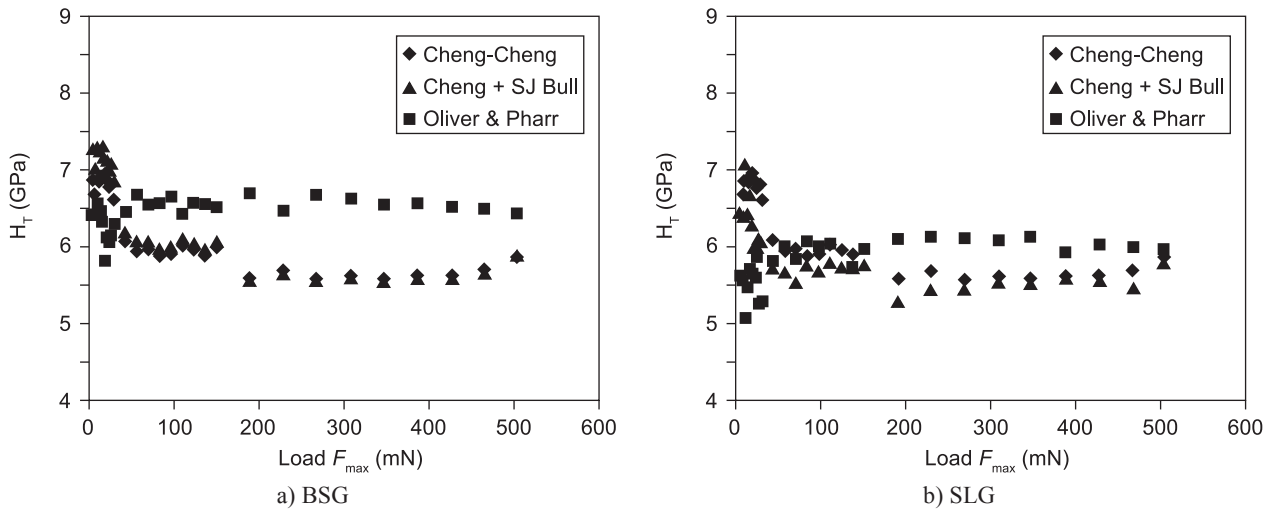


Figure 14. Comparison of variation of  $H_{IT}$  for BSG (a) and for SLG (b) obtained using energy methods (Cheng & Cheng, Chen & Bull) and Oliver & Pharr method [19,20,10].

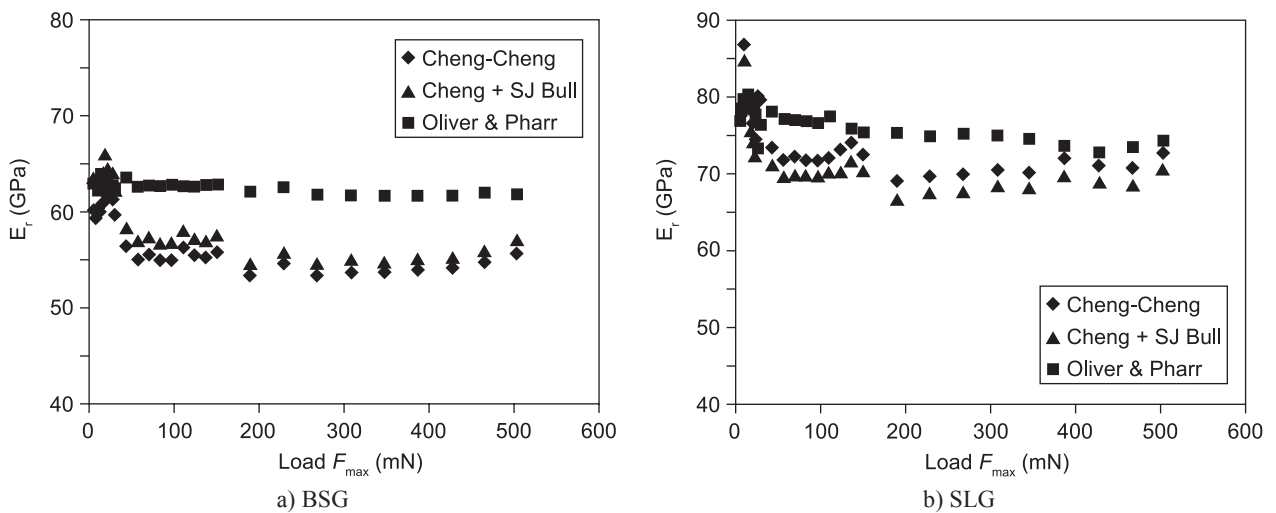


Figure 15. Comparison of variation of  $E_r$  for BSG (a) and for SLG (b) obtained using the energy methods (Cheng & Cheng, Chen & Bull) and Oliver & Pharr method [19,20,10].

### Comparison of Energy methods with Oliver and Pharr method

We also derived the two glasses properties ( $H_{IT}$  and  $E_r$ ) according to the nonlinear C&B energy method using the relations 13 and 14. The work hardening exponent  $n$  is null in our case and the others constants were taken as defined previously ( $\epsilon = 0.75$ ,  $\beta = 1.05$  and  $\alpha = 70.3^\circ$ ). Comparison of  $H_{IT}$  and  $E_r$  variations obtained by this method and those obtained by the linear (C&C) energy method revealed that they are close when  $\lambda = 5.3$ , particularly for the borosilicate glass, as shown in Figures 14 and 15. On the other hand, if we compare the reference properties (Table 2) with the results from the three methods at large depths, we can see that for the soda lime glass, they are close to those obtained by the energy methods whereas for the borosilicate glass, they are rather comparable to those obtained by O&P analysis.

### CONCLUSION

The comparison made between the displacement O&P method with those based on the indentation energy in determining the hardness and the elastic modulus of two glasses (soda lime glass and borosilicate glass) revealed the following observations:

The qualitative assessment of the properties  $H_{IT}$  and  $E_r$  obtained by O&P method depends on the calibrations performed for correcting different error sources. The loading frame compliance has an important influence on the evaluated properties. Its implementation in the analysis was verified by linear regression of the variation of the measured compliance in terms of peak load  $F_{max}$  and also by examining the constancy of the parameter  $F_{max}/S^2$  at large depths. The calibrated surface area obtained on fused silica led to a significant reduction in the apparent hardness ISE effect at small depths. The effects of the initial depth and the thermal drift are relatively less important in our case.

The application of the energy methods based on a linear relationship between  $W_e/W_t$  and  $H_{IT}/E_r$  showed a dependency of the constant  $\lambda$  on the tested material. For a value of  $\lambda = 5.3$ , the results obtained by this method are close to those obtained by the non linear method proposed by Chen and Bull. The comparison with the reference values obtained by conventional means revealed that these energy methods gave acceptable results for the soda lime glass. The borosilicate glass reference properties are rather close to those obtained by Oliver and Pharr method.

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